



## Impacts of leaks and gas accumulation on closed chamber methods for measuring methane and carbon dioxide fluxes from tree stems

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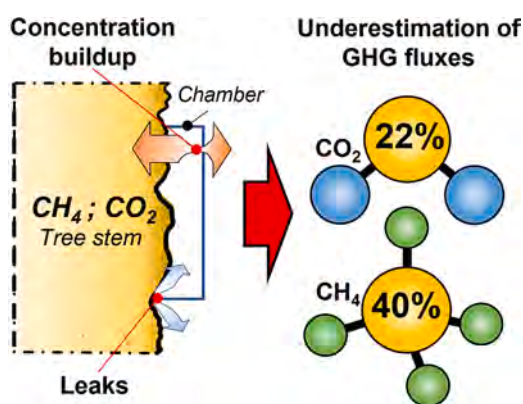
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### HIGHLIGHTS

- CH<sub>4</sub> and CO<sub>2</sub> fluxes from tree stems are significant and important to constrain.
- Stem flux measurements can be affected by leaks and concentration buildup effects.
- We studied these effects in six tree species from two contrasting forest ecosystems.
- These effects led to an underestimation of 40 % for CH<sub>4</sub> and 22 % for CO<sub>2</sub> fluxes.
- A simple method addresses both leaks and concentration buildup effects.

### GRAPHICAL ABSTRACT



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### ABSTRACT

Accurate measurements of methane (CH<sub>4</sub>) and carbon dioxide (CO<sub>2</sub>) fluxes from tree stems are important for understanding greenhouse gas emissions. Closed chamber methods are commonly employed for this purpose; however, leaks between the chamber and the atmosphere as well as gas accumulation, known as the concentration buildup effect, can impact flux measurements significantly. In this study, we investigated the impacts of concentration buildup and leaks on semi-rigid closed chamber methods. Field measurements were conducted on six tree species, including three species from a Mexican mangrove ecosystem and three species from a Magellanic sub-Antarctic forest. Systematic observations revealed significant leak flow rates, ranging from 0.00 to 465 L h<sup>-1</sup>, with a median value of 1.25 ± 75.67 L h<sup>-1</sup>. We tested the efficacy of using cement to reduce leaks, achieving a leak flow rate reduction of 46–98 % without complete elimination. Our study also demonstrates a clear and substantial impact of concentration buildup on CH<sub>4</sub> flux measurements, while CO<sub>2</sub> flux measurements were

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relatively less affected across all tree species studied. Our results show that the combined effects of leaks and concentration buildup can lead to an underestimation of CH<sub>4</sub> emissions by an average of  $40 \pm 20$  % and CO<sub>2</sub> emissions by  $22 \pm 22$  %, depending on the bark roughness. Based on these findings, we recall a straightforward yet effective method to minimize experimental errors associated with these phenomena, previously established, and reiterated in the current context, for calculating emissions that considers effects of leaks and concentration buildup, while eliminating the need for separate determinations of these phenomena. Overall, the results, combined with a literature review, suggest that our current estimates of GHG flux from tree stems are currently underestimated.

## 1. Introduction

In the context of climate change, understanding greenhouse gas (GHG) sources and sinks is of paramount importance. Forests are widely acknowledged as global sinks for methane (CH<sub>4</sub>; Saunois et al., 2020), where it is produced in the anaerobic deeper soil layer but oxidized in excess, within the aerobic upper layer. However, in flooded forests, like mangroves, water creates an oxygen barrier, turning forest soils into a net CH<sub>4</sub> source. Regardless of the specific forest type, it is well-established that trees play a pivotal role in the transport of CH<sub>4</sub> produced in forest soils (Barba et al., 2019; Covey and Megonigal, 2019). The internal structure of tree stems facilitates the transport of soil CH<sub>4</sub> through both bark aerenchyma and xylem tissues (Teskey et al., 2008; Bloemen et al., 2013; Vroom et al., 2022; Yáñez-Espinosa and Ángeles, 2022).

Tree-mediated emissions, therefore, represent a deviation from the standard CH<sub>4</sub> cycle in forest soils and can offset up to 46.5 % of the soil CH<sub>4</sub> uptake (Machacova et al., 2023). In fact, in the Amazon basin alone, approximately  $21.2 \pm 2.5$  megatons (Mt) of CH<sub>4</sub> are emitted by trees annually, accounting for 10–13 % of global wetland CH<sub>4</sub> emissions (Pangala et al., 2017). On a broader scale, Covey and Megonigal (2019) reported CH<sub>4</sub> emission rates from tree stems ranging from  $-0.014$  to  $6.5 \text{ g m}^{-2} \text{ d}^{-1}$ . Regarding carbon dioxide (CO<sub>2</sub>), while forests significantly contribute to carbon dioxide (CO<sub>2</sub>) sequestration, capturing about 11 gigatons (Gt) of CO<sub>2</sub> per year (Friedlingstein et al., 2022, emissions from tree stems contribute to 5–38 % of the total ecosystem respiration (Carnioli et al., 2016; Yang et al., 2016). Though the contribution of CO<sub>2</sub> emissions from tree stems to the global CO<sub>2</sub> budget is relatively small, it remains relevant for studying tree physiology (Teskey et al., 2017) throughout diurnal and/or seasonal cycles.

Measurement of GHG emissions from tree stems is therefore important, to constrain better our current estimates and to provide better inventories of GHG. Current methods for the determination of these emissions are mostly based on the use of chambers, temporarily fixed on tree stems in which the emitted gases are captured. Among several designs, semirigid chambers, as described by Siegenthaler et al. (2016), are commonly used due to their ease of installation in the field. Regardless of the chamber design, chambers can be operated with continuous (dynamic) or discrete (static) sampling and measurement of the gas concentration. When a detector is used for dynamic measurement, it is usually connected to the chamber through a closed loop, and thus referred to as closed chambers. In that case, fluxes are determined from the derivative of the gas concentration increase within the chamber. The same data processing strategy is generally used when discrete measurements are done.

Whatever the chamber design and the method used, flux measurements are challenging, mostly because they combine low flux magnitude, difficulties in hermetically coupling the chamber and the tree stems, and complex pneumatic behavior of the experimental setup, including the chamber and the pneumatic circuit connecting it to the gas analyzer. One of the challenges in flux measurements is to securely affix the chamber on tree stems avoiding leaks, i.e. gas exchange between the atmosphere and the internal volume of the chamber, that would bias the results. At least in some cases, the shape of tree stems or the natural bark roughness, which is highly variable among tree species, hinder hermetic

sealing, and a common strategy is the use of different sealants (e.g., moldable cement, potting clays, non-caustic silicone, flexible putty adhesive, silicone rubber) to improve chamber's hermeticity (Bréchet et al., 2021; Flanagan et al., 2021; Jeffrey et al., 2021; van Haren et al., 2021). However, the use of sealants is not an absolute guarantee of sealing, at least in the case of very rough bark presenting deep channels, fractures, or irregular shapes. Moreover, even if the chamber is perfectly sealed, leaks might appear in the pneumatic circuit or within the detector, as exemplified by Wilkinson et al. (2018). Leaks from different origins might therefore have a significant impact on the accuracy of emission measurements.

In addition to leaks, another potential challenge, is that during flux measurements an increase in concentration of the measured gases within the chamber is observed. This concentration buildup has the potential to reduce the gas partial pressure gradient between the tree stem and the atmosphere, which would otherwise be observed without a chamber. In turn, this phenomenon may result in underestimated fluxes, as suggested in aquatic ecosystems (Xiao et al., 2016), soils (Welles et al., 2001; Kutzbach et al., 2007; Juszczak, 2013) and observed in trees (Jeffrey et al., 2020; Bréchet et al., 2021; Kohl et al., 2021; van Haren et al., 2021).

Lastly, while this has not been systematically studied, our interpretation is that, in dynamic chambers, the experimental setup forms a complex pneumatic circuit, which can also pose challenges in interpreting the measured concentration profile. Indeed, the gas within the chamber, exhibiting relatively good mixing, is pumped to the analyzer through tubing where plug-flow is observed. Subsequently, the gas passes through the sensor cavity of the analyzer, with its own pneumatic behavior, before returning to the chamber, in the case of closed chambers. In consequence, during the starting period of the measurements, these phenomena can lead to delays, mixing effects, and potential concentration gradients, which hinder the interpretation of the gas concentration profiles observed. These effects have been previously identified and referred to as “dead band” by Siegenthaler et al. (2016) and it is a common practice to discard data from the beginning of the measurements when continuous gas concentration is recorded (Pitz and Megonigal, 2017; van Haren et al., 2021; Epron et al., 2022; Fraser-McDonald et al., 2022).

Overall, the impact of leaks and gas concentration buildup, along with the accurate interpretation of the pneumatic behavior in flux measurement setups, require a systematic and experimental analysis. This is the main objective of the present study. To achieve this goal, we established a mass balance of closed chambers and developed a method to quantify leaks and to analyze the effect of concentration buildup within the chamber headspace, with the hypothesis that by considering and quantifying leaks and analyzing the effect of concentration buildup, the accuracy of greenhouse gas (GHG) flux measurements from tree stems will be improved. The developed method was applied to six different tree species with distinct characteristics, from two contrasting ecosystems: a tropical flooded mangrove in Mexico and a subantarctic non-flooded forest in southern Chile. Through our findings, we quantified the potential errors associated with measuring GHG fluxes from tree stems and discussed various experimental procedures commonly documented in the literature.

## 2. Materials and methods

### 2.1. Chamber mass balance

A detailed mass balance of the dynamic closed chamber and the corresponding equations are presented in Supplementary Material (Section S1.1). Briefly, the time derivative of the gas concentration within the chamber ( $C_C$ ) can be summarized into fluxes from the tree stem and leaks between the atmosphere and the chamber headspace or the pneumatic circuit (Fig. 1).

$$\frac{dC_C}{dt} = F \cdot \frac{A_C}{V_T} + k_L \cdot (C_A - C_C) \quad (1)$$

where  $F$  is the observed flux;  $A_C$  is the chamber area in contact with the tree stem;  $V_T$  is the total volume of the chamber, the detector and the pneumatic circuit;  $k_L$  is the leak constant ( $h^{-1}$ ), which represents the frequency at which the gaseous content of the chamber and the pneumatic circuit is replaced by external air due to leakage; and  $C_A$  is the atmospheric gas concentration, observed during the measurements. It is worth noting that  $k_L$  is the ration between the leaks flow rate ( $Q_L$ ) and  $V_T$ . Thus, leaks can be expressed, alternatively, by  $k_L$ , or by  $Q_L$ .

In Eq. (1),  $F$  is the flux measured when the chamber is positioned on the tree stem. Thus, the surface of the bark is exposed to  $C_C$ , potentially different from the atmospheric concentration  $C_A$ . Hence, the actual flux ( $F^*$ ), naturally occurring when the bark is exposed to  $C_A$  (no chamber) might be different from  $F$ . Since mass transfer is a lineal function of the concentration gradient, we can estimate that the measured flux is a function of the ration between the concentration gradient with and without a chamber:

$$F = F^* \cdot \frac{(C_i - C_C)}{(C_i - C_A)} \quad (2)$$

where  $(C_i - C_C)$  and  $(C_i - C_A)$  are the concentration gradient observed during the chamber measurement and the gradient that would be otherwise observed in the absence of a chamber, respectively. Hereafter, Eq. (2) will be referred to as “the concentration buildup effect”, which bears similarity to equations commonly used in the study of aquatic ecosystems (Xiao et al., 2016) and wetlands (Kutzbach et al., 2007; Juszczak, 2013). Therefore, considering this concentration buildup effect, Eq. (1) becomes:

$$\frac{dC_C}{dt} = F^* \cdot \frac{(C_i - C_C)}{(C_i - C_A)} \cdot \frac{A_C}{V_T} + k_L \cdot (C_A - C_C) \quad (3)$$

It is worth noting that, in an ideal case, at the onset of the measurement, just after chamber installation,  $C_C$  might be equal to  $C_A$ , and

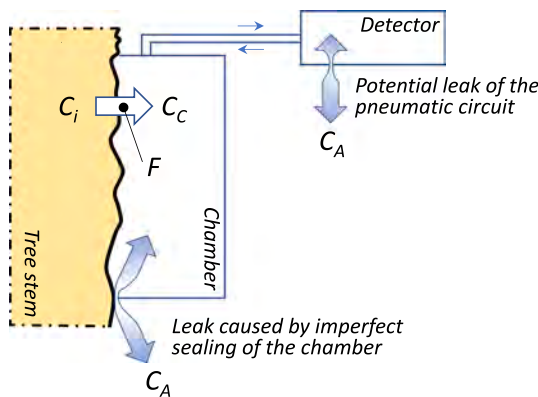


Fig. 1. Conceptual scheme of the chamber mass balance. Where  $F$  is the flux from the tree stem, and  $C_C$ ,  $C_A$  and  $C_i$  are the gas concentration within the chamber, the atmospheric gas concentration, and the tree internal gas concentration, respectively.

Eq. (3) is greatly simplified, as follows:

$$\frac{dC_C}{dt} = F^* \cdot \frac{A_C}{V_T} \quad (4)$$

Eq. (4) holds significant importance as it suggests that accurately measuring flux, free from the effects of leaks and concentration buildup, can be achieved by analyzing the initial slope of gas concentration within the chamber (Eq. (4)). However, in practice, measuring this initial slope, when  $C_C$  equals  $C_A$ , presents challenges for two primary reasons. Firstly, at the start of the measurement, there may be differences in gas concentration among the chamber, the pneumatic circuit, and the detector, leading to a transient period as the system begins to operate. This is the main reason why it is a common practice to disregard data from the beginning of the measurements when calculating fluxes using continuous sampling (Siegenthaler et al., 2016; Pitz and Megonigal, 2017). Secondly, during the installation of the chamber, gases may already accumulate within it or be influenced by the proximity of operators (often notable with  $CO_2$  expelled by the operators), resulting in  $C_C$  being higher than  $C_A$ . Consequently, in real-case scenarios, even at the onset of the measurements,  $C_C$  often exceeds  $C_A$  and the observed concentration increase is influenced by multiple mechanisms beyond just the flux. Under these conditions, Eq. (4) becomes inappropriate, and ideally, Eq. (3) should be employed to accurately quantify the flux.

In Eq. (3), if the time interval  $dt$  is sufficiently small, we can estimate the change in  $C_C$  over time ( $C_{C,t}$ ), after the chamber has been installed and operated, by adding the rate of change ( $dC_C/dt$ ) multiplied by the small-time interval ( $\Delta t$ ) to the previous value of  $C_C$  ( $C_{C,t-1}$ ). Mathematically, this approximation can be expressed as follows:

$$C_{C,t} = C_{C,t-1} + \left( F^* \cdot \frac{(C_i - C_{C,t-1})}{(C_i - C_A)} \cdot \frac{A_C}{V_T} + k_L \cdot (C_A - C_{C,t-1}) \right) \cdot \Delta t \quad (5)$$

In Eq. (5), the geometry of the chamber ( $A_C$ ,  $V_T$ ) is known, and  $C_A$  can be readily measured in the field. Therefore, if  $k_L$  is known, by calibrating Eq. (5) to the experimental  $C_C$  time series, it becomes possible to determine  $F^*$  and  $C_i$ , effectively accounting for the effects of leaks and concentration buildup.

For the determination of  $k_L$ , we developed a method, that is described in detail in the Supplementary Material (Section S1.2). Briefly, the method involves inducing a transient state by injecting a pulse of  $CH_4$  into the chamber, artificially creating a relatively high concentration of  $CH_4$ . Following this injection, any leakage would result in an asymptotic decrease in the  $CH_4$  concentration, until the system reaches a new steady state. Assuming a well-mixed behavior of the gas phase within the chamber, this asymptotic decrease could be described by Eq. (6), where  $C_{C,0}$  and  $C_{C,f}$  are the  $CH_4$  chamber concentration at time zero (shortly after pulse injection) and after the chamber reaches steady state (details provided in Section S1.2).

$$C_{C,t} = C_{C,0} + (C_{C,0} - C_{C,f}) \cdot (1 - \exp(-k_L \cdot t)) \quad (6)$$

By calibrating Eq. (6) to the experimental  $C_C$  time series, it is possible to determine  $C_{C,f}$  and  $k_L$ , with the latter being the parameter of primary interest. However, it is important to note that for this purpose,  $C_{C,0}$  must be sufficiently high, typically a few hundreds of ppm, to ensure a subsequent decrease in  $C_{C,t}$  over time. Furthermore, in the Supplementary Material, we demonstrate that while Eq. (6) is correct, it does not fully capture the intricate pneumatic behavior observed in certain cases (Fig. S2). Although we have mathematically described this complex behavior, it falls outside the primary scope of the present study, which is centered around evaluating the impact of leaks on emissions. Therefore, we have included these findings in the Supplementary Material, Section S1, provided for those readers who are interested in exploring this aspect further.

## 2.2. Fluxes, leaks and $C_i$ determination

In the present work, fluxes of  $\text{CH}_4$  and  $\text{CO}_2$  ( $F_{\text{CH}_4}^*$  and  $F_{\text{CO}_2}^*$ , respectively), were determined according to the following 4-step protocol (materials being described in Section 2.3). Step 1, semi-rigid chambers were tidily fixed on the tree stem using 3 to 4 nylon straps. Step 2, the air content of the chamber where  $\text{CH}_4$  or  $\text{CO}_2$  potentially accumulated during chamber installation, was replaced by fresh ambient air over 2 min, with a portable external air pump (Flextailgear Tiny Pump, Mexico). Step 3, the chamber was immediately connected, in a closed loop, to a laser ultraportable greenhouse analyzer (i.e. UGGA, model 915-0011-1000, Los Gatos Research, ABB, USA). Step 4, the  $\text{CH}_4$  and  $\text{CO}_2$  concentration within the chamber was measured for 10–15 min. Step 5,  $k_L$  was determined from the injection of 1 mL of  $\text{CH}_4$  into the chamber. Pulse injection and  $k_L$  determination was done after flux measurement, to avoid any potential effect of increasing artificially  $\text{CH}_4$  concentration to much higher levels than standardly observed during flux determination. Specifically,  $k_L$  was determined from Eq. (6) calibration (details provided in Section S1.2). Step 6,  $C_C$  dataset measured during step 4, was used to calibrate Eq. (5), which was adjusted to the experimental data, using  $C_i$  and  $F^*$  as adjustment parameters, and where  $k_L$  was determined during step 5.

## 2.3. Experimental setup

Depending on the size of the tree, we used four different sizes of semi-rigid chambers (#1,  $0.50 \times 0.30$ ; #2,  $0.30 \times 0.24$ ; #3  $0.24 \times 0.18$ ; #4,  $0.15 \times 0.10$  m), all of them composed of a 0.7 mm thick impermeable polyethylene terephthalate (PET) plastic sheet glued along its entire perimeter to a  $20 \times 30$  mm closed cell neoprene foam (Seals+Direct Ltd., Hampshire, UK), as described in Siegenthaler et al. (2016). The distance between the tree stem and the PET plastic sheet, which determined the chamber volume, was 20 mm. The chambers were connected in a closed loop to the UGGA, with two 6 mm external diameter (4 mm internal diameter) flexible polyurethane tubing (PUN-6X1-DUO-BS, Festo, Mexico). To reduce the volume of the pneumatic circuit and improve the recirculation rate between the chamber and the detector, the length of this tubing was as short as possible (about 2 m) in dry ecosystems, but of up to 12 m long in the mangrove ecosystem, to ensure that the detector was placed in a dry and safe place. It should be noted that, in contrast to the relatively heavy UGGA model used in the present work (16 kg plus batteries), the current trend is towards lightweight and backpack portable GHG detectors, making their deployment in close vicinity of the chamber easier. The volume of the semi-rigid chambers used is an important parameter, which is required for flux determination (Eq. (3)). Compared to rigid chamber, which have an approximately constant volume, the volume of semi-rigid chambers depends on the deformation of the chamber around the tree stem. The volume of the chambers was determined according to Siegenthaler et al. (2016), and ranged 0.22–2.9 L. On the contrary to several previous reports, in the present work, the relatively low thickness of the chambers used impeded the use of an internal fan to homogenize the chamber headspace. However, the relatively small chamber volume and the UGGA flow rate ( $1.2 \text{ L min}^{-1}$ ) ensured a relatively good mixing of the chamber volume. In some experiments, the impact of moldable cements (Play-Doh, purchased in toy stores, México) was used as leak repressor. In each case, flux and leak measurements were done according to the standard method described above.

## 2.4. Sites description and campaigns

The method was applied to two contrasting forest ecosystems; a mangrove forest located in Celestun, in the Yucatan Peninsula (Mexico; 20.855,  $-90.374$ ) and a sub-Antarctic forest located in Puerto Williams, in the Navarino Island (Chile;  $-54.949$ ,  $-67.660$ ). Celestun is characterized by a hot and semi-arid climate with a marked rainfall season

(Orellana et al., 2010), an annual mean rainfall of 746.7 mm (mean monthly range from 8.2 to 144.0 mm) and an annual mean temperature of  $26.4 \text{ }^\circ\text{C}$  (mean monthly range from  $23.2$  to  $28.5 \text{ }^\circ\text{C}$ ) (CONAGUA, 2023). In turn, Puerto Williams has a tundra and oceanic climate with an annual mean rainfall of 822 mm (mean monthly range from 51 to 131 mm) and an annual mean temperature of  $6.1 \text{ }^\circ\text{C}$  (mean monthly range from  $-2.8$  to  $7.2 \text{ }^\circ\text{C}$ ) (Aguirre et al., 2021; DMC, 2023). In the mangrove ecosystem, we measured flux and leaks on the stem of the tree dominant species; “red mangrove” (Rm; *Rhizophora mangle*), “black mangrove” (Ag; *Avicennia germinans*), and “white mangrove” (Lr; *Laguncularia racemosa*). In Navarino Island, we also characterized the three dominant species; “low deciduous beech” (Na; *Nothofagus antarctica*), “high deciduous beech” (Np; *Nothofagus pumilio*), and “evergreen beech” (Nb; *Nothofagus betuloides*). It is worth mentioning that, in each ecosystem, the bark of the trees was drastically different among species (Figs. S3 and S4), offering a good model to evaluate the impact of leaks. For each species, nine to thirteen trees were selected with a range of diameter at breast height (DBH) from 3.8 to 56.6 cm. Then flux, leaks and bark roughness were determined in each tree at two stem heights, at 0.25 m and 1.4 m (up to 26 measurements by species). In the case of Rm, characterized by stilt roots, all measurements were made at 0.25 m above the last stilt root. According to the stem diameter, chambers #1 and #2 were used for larger diameters mostly at a tree height of 0.25 m, chambers #2 or #3 were used for medium diameters at both tree heights, and chambers #3 and #4 were used for smaller diameters (see Table S2 for further details). The diameter of the trees was measured using a sewing measuring tape (Singer 50,003 ProSeries) which was calibrated with a caliper (Mitutoyo 530–101). As the objective of the present work was to quantify leaks and the effects of concentration buildup, rather than focusing on reporting emissions, we determined fluxes during the daytime only, typically from 10 am to 4 pm. The field campaign in México took place from September 6th to September 18th, 2021, while the Navarino campaign took place from February 20th to March 3rd, 2022. During these campaigns, the mean temperature was  $30.5 \pm 2.0 \text{ }^\circ\text{C}$  and  $8.9 \pm 3.1 \text{ }^\circ\text{C}$  for Celestun and Puerto Williams, respectively. While the mean wind speed was  $1.2 \pm 0.5 \text{ km h}^{-1}$  and  $6.8 \pm 4.5 \text{ km h}^{-1}$  for Celestun and Puerto Williams, respectively. Weather data were obtained from public data of nearby stations, namely, from the Celestun weather station (#31040; CONAGUA, 2023) located 2800 m west of the mangrove location, and from Puerto Williams airport station (#550001; DMC, 2023), located 3200 m northeast from the forest where measurements were taken.

## 2.5. Bark roughness determination

In order to correlate leaks to bark roughness, which is on first approximation a governing parameter of leaks, we developed our own method, founded on our field experience. The method is based on a 0.2 m contour gauge (6" Contour Gauge, General Tools & Instruments, Mexico) that was firmly pressed against the tree bark at the same height where chamber was placed. We defined bark roughness ( $R$ ; Fig. 2) as the relative difference between the length of the bark contour ( $L'$ ) and its equivalent smooth length ( $L$ ; Eq. (7)). In each case,  $L'$  was determined from a scaled photograph of the gauge taken in the field, and analyzed with ImageJ (Schneider et al., 2012).

$$R = \frac{(L' - L)}{L} \bullet 100 \quad (7)$$

Since the bark geometry of the tree stems differs along both axes, we distinguished axial and radial roughness, as follows. Axial roughness ( $R_{A\%}$ ) was defined as illustrated in Fig. 2A and Eq. (8):

$$R_{A\%} = \frac{(L'_A - L_A)}{L_A} \bullet 100 \quad (8)$$

In turn, radial roughness ( $R_{R\%}$ ) was defined as the perimeter section

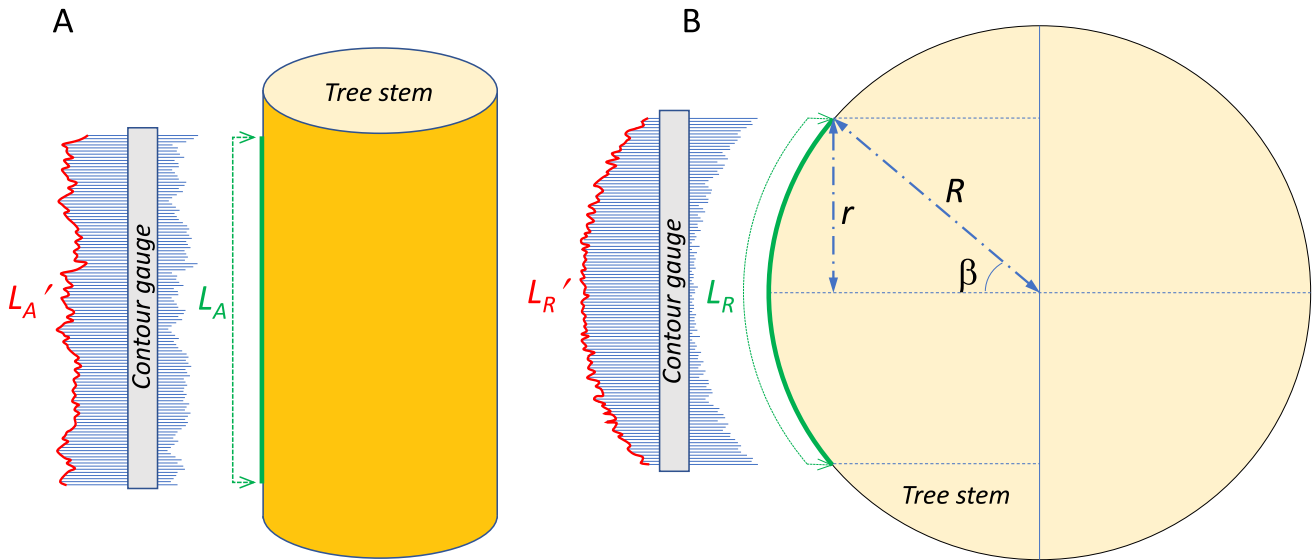


Fig. 2. Conceptual scheme of bark roughness for both axial (A) and radial (B) determination in a longitudinal and cross sections, respectively.

of the bark contour ( $L'_R$ ) and the section of a smooth circle of the same tree radius ( $L_R$ ), as shown in Eq. (9) and Fig. 2B.

$$R_{R\%} = \frac{(L'_R - L_R)}{L_R} \cdot 100 \quad (9)$$

In Eq. (9),  $L_R$  can be determined, as follows, where angles are expressed in radians and where  $r$  is half the length of the contour gauge;

$$L_R = \frac{2\beta}{2\pi} \cdot 2\pi R = 2\beta R \quad (10)$$

$$\sin(\beta) = \frac{r}{R}; \beta = \arcsin\left(\frac{r}{R}\right) \quad (11)$$

$$L_R = 2 \cdot R \cdot \arcsin\left(\frac{r}{R}\right) \quad (12)$$

For each measurement,  $R_{A\%}$  was taken from both lateral sides of the chamber and  $R_{R\%}$  was taken from upper and lower sides of the chamber. For each position, a single mean bark roughness percentage ( $R_{\%}$ ) was determined, using a weighted average (Eq. (13)), where  $H_C$  and  $W_C$  are the height and width of the chamber used.

$$R_{\%} = \frac{R_{A\%} \cdot H_C + R_{R\%} \cdot W_C}{H_C + W_C} \quad (13)$$

### 2.6. Data treatment and statistics

Eqs. (5) and (S8) were calibrated to experimental data using a Generalized Reduced Gradient (GRG) non-linear tool and minimizing the root-mean-square error (RMSE) between experimental data and models. Normality (Shapiro-Wilk) and homoscedasticity (Levene) were tested prior to statistical analysis. Then, two-way ANOVA were performed. Most variables had a positive skew, which were log-transformed to achieve normality. If normality was not achieved, then a Kruskal-Wallis test was applied instead. Statistical differences among tree species, chamber designs, and ecosystems, were determined according to Tukey's post-hoc tests. All statistical analysis were performed in R language version 2023.03.1.446 (Posit Team, 2023) or with Origin(Pro) software (OriginLab Corporation, 2016). In the present work, results are usually presented in term of the median  $\pm$  one standard deviation, unless specified. Boxplots in Section 3 and Supplementary Material display the first quartile (Q1), median (Q2), and third quartile (Q3) as boxes, with whiskers extending to the minimum and maximum values. Outliers are represented as individual data points beyond the whiskers.

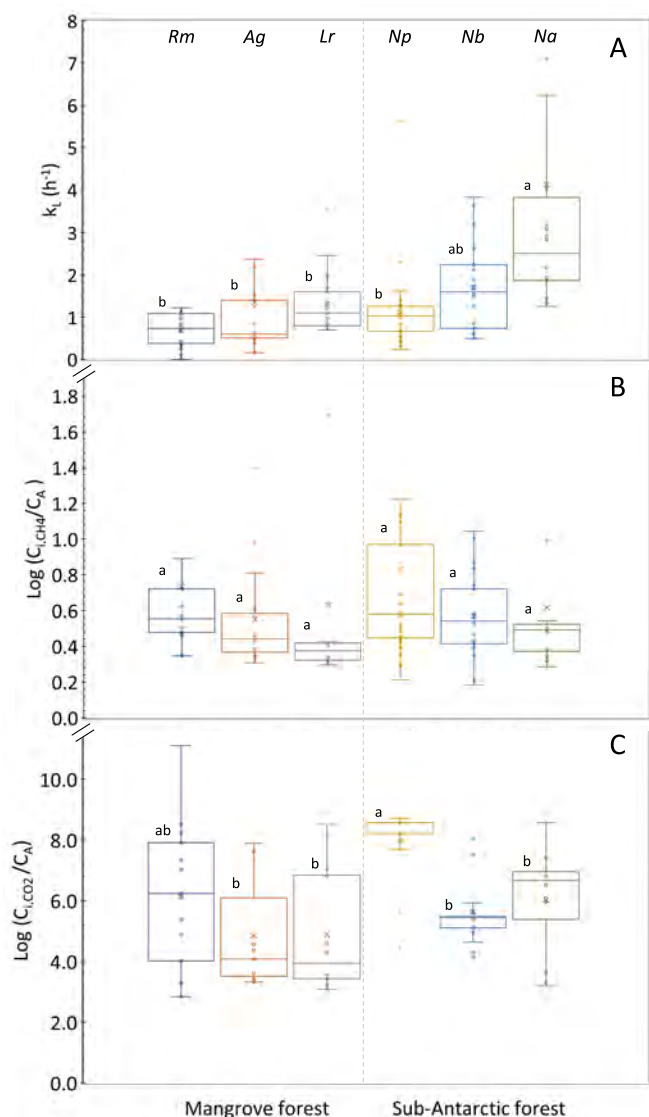
## 3. Results and discussion

### 3.1. Leaks characterization

Fluxes and leaks were determined 153 times in different trees and under different configurations. Overall  $k_L$  ranged 0.00–284  $h^{-1}$ , with a median of  $1.25 \pm 46.9 h^{-1}$ , which corresponded to leaks flow rates ( $Q_L$ ) from 0.00 to 465  $L h^{-1}$ , with a median of  $1.25 \pm 75.67 L h^{-1}$ . Thus, the chambers were far from being well sealed, and significant gas exchanges occurred between the chamber headspaces and the atmosphere. It is worth noting that the leak flow rates reported here are not exclusively caused by the chamber sealing, but include all components of the measurement setup, including the UGGA leak flow rate of  $0.27 \pm 0.05 L h^{-1}$ , reported in Section S1.3. Compared to the total leak flow rate, the UGGA's leak represented, on average, about 8.8 % of the total leaks.

No significant difference of leaks was observed among tree species, except Na, which were significantly higher than those observed with Rm, Ag, Lr, and Np ( $P < 0.05$ ; Fig. 3A). The notable difference in leaks between Na and the other tested species was most certainly due to bark roughness, as  $R_{\%}$  of Na ( $51.1 \pm 16.8 \%$ ) was significantly higher than all other species ( $P < 0.05$ ; Fig. S4). It is important to mention that there were no differences between the roughness measured at two stem heights within each species (Table S1). To further characterize possible correlations between bark roughness and leaks,  $R_{\%}$  of the measured trees was compared to  $k_L$ . A positive exponential correlation was observed between  $R_{\%}$  and  $k_L$  ( $P < 0.05$ ; Fig. S5). On the contrary, no significant difference was observed among  $k_L$  measured in the same trees, but with different chambers ( $p = 0.66$ ), which is not surprising because  $k_L$  is a specific parameter, i.e. where  $Q_L$  is proportional to  $V_T$ , (Fig. S6).

Leaks between the chamber and the atmosphere might be dependent on the weather conditions, and particularly on the wind speed, which might alter the gas exchange between the chamber and the atmosphere, leading to flux measurement errors or artifacts (Bain et al., 2005; Maier et al., 2019; Jiang et al., 2023). To check for the latter, in another set of experiments, 6 replicates of  $k_L$ , were determined over two hours, in two trees; one presenting a moderate roughness (Np;  $R_{\%} = 4.10$ ) and one presenting a high roughness (Na;  $R_{\%} = 88.58$ ). This test was done on a day with variable winds, i.e.,  $18.7 \pm 6.0 km h^{-1}$  with gusts at  $39 km h^{-1}$ . In the case of Np, the mean  $k_L$  was equal to  $0.86 \pm 0.02 h^{-1}$ , with a coefficient of variation of 2.0 %, which suggest that the variable wind had little effects on leak determination. This was confirmed by a generally good adjustment of the leak model (Eq. (S8)) to the experimental data, with a mean  $R^2$  of  $0.999 \pm 0.001$ . On the contrary, in the



**Fig. 3.** Leaks, expressed as  $k_L$ , among the different tree species (A); ratio between  $C_i$  and  $C_A$ , for  $CH_4$  (B) and  $CO_2$  (C).

case of Na, the mean  $k_L$  was  $16.8 \pm 6.2 \text{ h}^{-1}$  (coefficient of variation of 37%), and with a poorer model fitting ( $R^2$  of  $0.951 \pm 0.109$ ). Therefore, the wind conditions had a major impact on leak detection in the tree with high roughness, and it would be advisable, in these cases, to determine fluxes during stable and moderate wind conditions or to use cement to reduce  $k_L$ .

### 3.2. Impact of moldable cement

The most common strategy to avoid or reduce leaks is using cement. To evaluate the impact of cement, another set of measurements was done with and without moldable cement on trees where significant leaks were previously detected. Overall, the use of cement reduced the leak flow rates by 46–98%, with a mean of  $92 \pm 20\%$  ( $n = 20$ ), which corresponded to a reduction of  $k_L$  by  $92 \pm 27\%$ . Thus, leaks were not completely avoided by cement, but they were reduced to a large extent (data not shown).

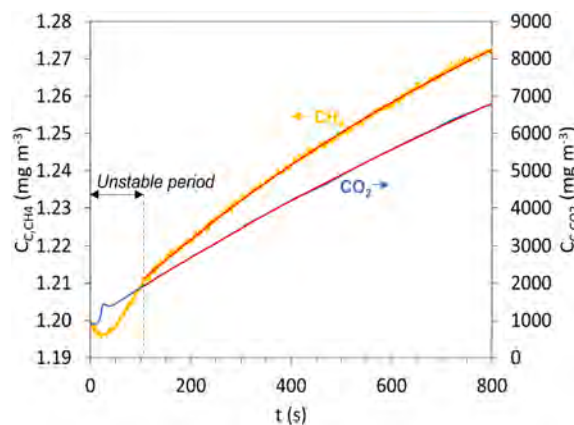
### 3.3. Fluxes and $C_i$ determination

As expected from Eq. (5), an asymptotic trend of  $C_C$  was almost

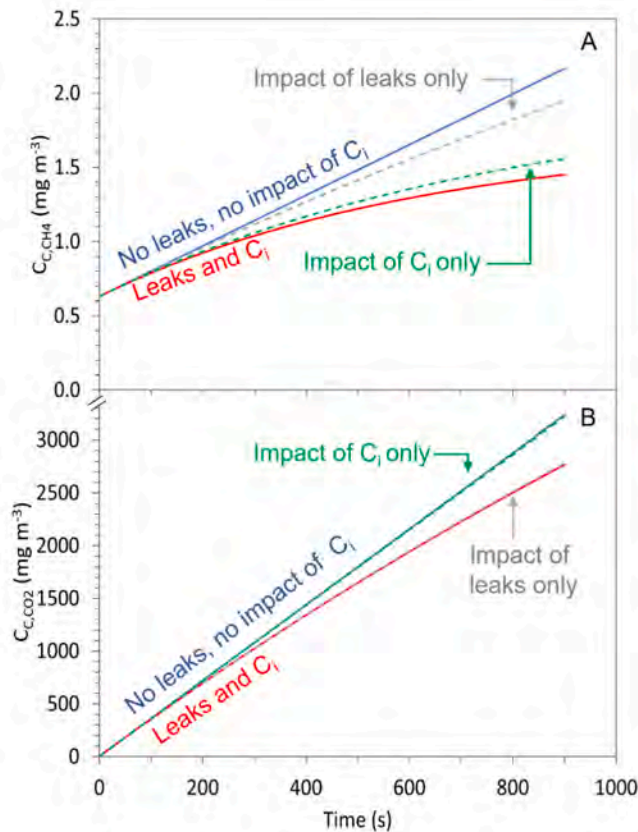
systematically observed, more clearly marked for  $C_{C,CH_4}$  than for  $C_{C,CO_2}$ , as exemplified on Fig. 4. Notably, at the onset of the measurements, and despite the chamber ventilation performed to allow for equilibration of the experimental setup, an instability that lasted 2–3 min was standardly observed (Fig. 4). Therefore, calibration of Eq. (5) was done discarding this initial period. Overall,  $F_{CH_4}^*$  ranged 0–2.05  $\text{mg m}^{-2} \text{h}^{-1}$  with a mean of  $0.17 \pm 0.29 \text{ mg m}^{-2} \text{h}^{-1}$ . These emissions are in accordance with the literature, i.e.  $-0.06$  to  $271.68 \text{ mg m}^{-2} \text{h}^{-1}$  (Covey and Megonigal, 2019), with an estimated mean from these reported values of  $22.16 \pm 53.40 \text{ mg m}^{-2} \text{h}^{-1}$ . Regarding  $CO_2$ ,  $F_{CO_2}^*$  ranged 32–2851  $\text{mg m}^{-2} \text{h}^{-1}$  with a mean of  $358 \pm 378 \text{ mg m}^{-2} \text{h}^{-1}$ . This range of  $F_{CO_2}$  range is also in accordance with the literature, i.e. 0.33 to  $406.05 \text{ mg m}^{-2} \text{h}^{-1}$  (Campioli et al., 2016; Yang et al., 2016; Salom3n et al., 2017), with an estimated mean from these reported values of  $112.55 \pm 100.38 \text{ mg m}^{-2} \text{h}^{-1}$ . In the case of mangrove trees, the mean stem  $CH_4$  flux was  $0.10 \pm 0.17 \text{ mg m}^{-2} \text{h}^{-1}$ , which is slightly greater than the mean stem  $CH_4$  flux of  $0.04 \pm 0.05 \text{ mg m}^{-2} \text{h}^{-1}$  previously reported for mangroves of the same genera (He et al., 2019; Jeffrey et al., 2019; Zhang et al., 2019; Dušek et al., 2021; Gao et al., 2021). Concerning the *Nothofagus* species measured on Navarino Island, the mean stem  $CH_4$  flux was  $0.06 \pm 0.15 \text{ mg m}^{-2} \text{h}^{-1}$ , with no previous reports yet.

By model calibration, we observed that  $C_{i,CH_4}$  was varying largely and was relatively close to the atmospheric concentration. To reflect the latter more evidently, we expressed the closeness of  $C_i$  to  $C_A$  by the ratio between them. The  $C_{i,CH_4}/C_{A,CH_4}$  ratio was  $3.2 \pm 11.5$  (median  $\pm$  one standard deviation; Fig. 3B). The latter indicate that the  $CH_4$  fluxes measured were highly sensitive to the concentration buildup phenomenon, and that this observed behavior was shared by the six tested tree species, with no significant difference (Fig. 3B). The same approach with  $CO_2$  provided a very distinct behavior, with a median  $C_{i,CO_2}/C_{A,CO_2}$  ratio of  $337 \times 10^3 \pm 19 \times 10^3$  (Fig. 3C), suggesting that, in most cases, no significant effect of  $CO_2$  concentration buildup within the chamber was observed. Although Np was characterized by a  $C_{i,CO_2}$  significantly higher than those of the other tree species, all species exhibited a  $C_{i,CO_2}/C_{A,CO_2}$  several order of magnitude above  $C_{i,CH_4}/C_{A,CH_4}$  (Fig. 3B and C).

To further visualize the segregated impact of leaks and concentration buildup during flux measurements, Fig. 5 shows the  $CH_4$  and  $CO_2$  concentration profiles in a chamber, established numerically, from the mean results obtained in the present work, i.e. median  $A_C$ ,  $V_T$ ,  $F^*$ ,  $k_L$  and  $C_i$ . It can be observed that, compared to the linear increase of  $C_C$  expected in an ideal system, the presence of leaks and the effects of concentration buildup, whether separately or combined, resulted in a similar curvature in the  $C_C$  trends. However, the contribution of both effects is drastically different among  $CH_4$  and  $CO_2$  trends. In the case of  $CH_4$ , most of the curvature effect is caused by the concentration buildup, while in the case of  $CO_2$ , the effect of  $C_i$  appears insignificant. Overall, 76% of the deviation from linear  $C_{C,CH_4}$  increase was explained by the



**Fig. 4.** Example of  $C_C$  measurement and model fitting (Eq. (5); red continuous lines).

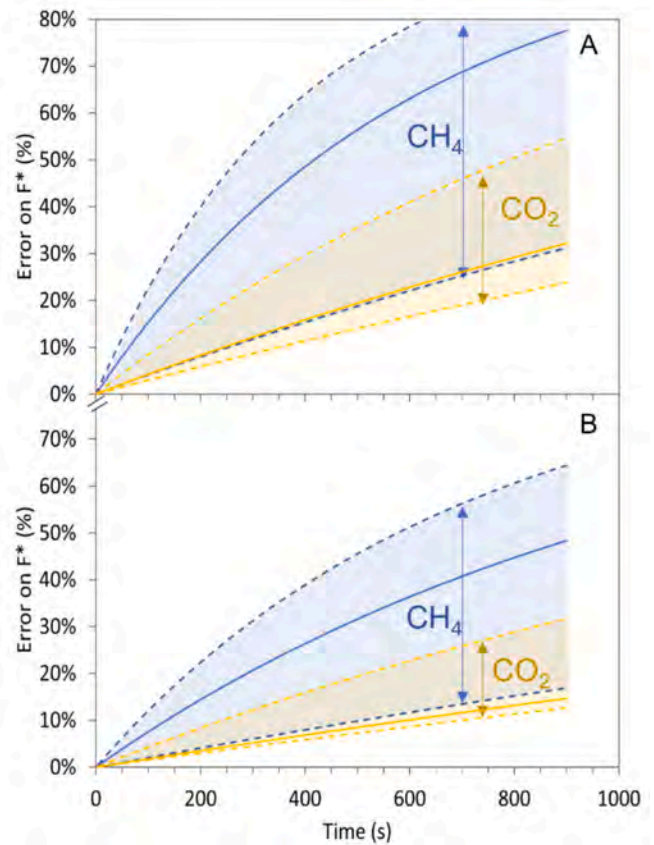


**Fig. 5.** Theoretical profiles of  $C_{C,CH_4}$  (A) and  $C_{C,CO_2}$  (B) that are expected during chamber deployment, from median  $A_C$ ,  $V_T$ ,  $F^*$ ,  $k_L$  and  $C_i$ , without considering leaks and the effect of  $C_i$  (ideal and linear profile), considering the effect of leaks only (median  $k_L$ ), considering only the concentration buildup (median  $C_i$ ), and considering both median  $k_L$  and  $C_i$ .

effect of concentration buildup, and 24 % was explained by leaks. The same exercise for  $CO_2$ , indicated that 3 % of the deviation to the linear  $C_{C,CO_2}$  increase was explained by the effect of concentration buildup, and 97 % was explained by leaks ( $k_L$ ). Therefore, the relative weight of both mechanisms is drastically different among gases.

### 3.4. Errors in $F$ determination

At the beginning of the measurements, when the chamber has been properly ventilated with atmospheric air and if  $C_C$  is equal to  $C_A$ , the initial slope of  $C_C$  accurately reflects the actual flux, independent of  $k_L$  and  $C_i$ . Therefore, one common error in determining  $F$  would be to assume that the initial linear slope is correct, because independent of leaks and concentration buildup effects. This would be an error, because disregarding potential gas accumulation during chamber installation or any other influences in the early stages of measurement. To estimate this potential error, we conducted a numerical analysis to determine the range of errors in flux determination, considering a delay between the start of chamber operation ( $t = 0$ ,  $C_C$  equals  $C_A$ ) and the actual acquisition of  $C_C$  data. In this exercise, we focused exclusively on the effect of gas accumulation resulting from tree emissions during chamber deployment, as other potential effects are stochastic and challenging to quantify. Based on the mean, second, and third quartiles of  $k_L$ ,  $F^*$ , and  $C_i$  complete dataset, we compared the slope of  $C_C$  observed after a delay time with the theoretical slope at time zero (Eq. (3)). Fig. 6A shows that even a short delay of three minutes can result in an underestimation ranging 7–37 % in  $F_{CH_4}$  determination and 5–15 % in  $F_{CO_2}$  determination.



**Fig. 6.** Range of error theoretically committed when discarding delay in starting measurements of  $C_C$  (A) and as a function of the duration of the measurements (B), for  $CH_4$  (light blue area) and for  $CO_2$  (light yellow area). Continuous lines show the error considering the median  $k_L$ ,  $F^*$ , and  $C_i$  data observed in the present work, while dashed lines shows the limit between the second and third quartile of these parameters.

A second, more obvious error in determining  $F^*$  would be to neglect the effects of leaks and concentration buildup and to assume that the slope between the initial  $C_C$  value and  $C_C$  at any time adequately reflects the flux. This error would occur when using a discrete sampling strategy or employing linear regression for dynamic  $C_C$  measurements. To assess the magnitude of this error, in this case too, we calculated the error range using the second and third quartiles of our complete dataset of  $k_L$ ,  $F^*$ , and  $C_i$ . The results (Fig. 6B) are that, for instance, a short 3-min measurement would result in an underestimation ranging from 4 to 20 % for  $F_{CH_4}$  and from 3 to 8 % for  $F_{CO_2}$  determination. It is important to note that the error presented in Fig. 6B does not account for the delay-induced error shown in Fig. 6A, which further emphasizes the magnitude of errors, when combined. Moreover, Fig. 6B highlights that the errors resulting from discrete measurements over extended durations reach magnitudes that potentially exceed an acceptable threshold. For instance, a 15-min measurement may result in a range of underestimation from 17 % to 64 % for  $F_{CH_4}^*$ , not to mention the potential errors when there are several hours between measurements, as sometimes reported.

To estimate the experimental error, combining the errors caused by delay and experimental duration, we determined fluxes in each  $C_C$  profile measured in the present work, discarding effects of leaks and concentration buildup, and considering in each case only the first three minutes of measurement, after the initial instable period. Under these conditions, the error committed on  $F_{CH_4}^*$  was an underestimation by 1.3–92 % with a median of  $40 \pm 20$  %, and an underestimation by 0–91 % with a median of  $22 \pm 22$  %, regarding  $F_{CO_2}^*$ , similar to the 33 %

underestimation in CH<sub>4</sub> fluxes from soils found by Pihlatie et al. (2013) in static chambers.

### 3.5. Practical considerations

The chamber model and the results obtained indicate that leaks and concentration buildup during flux measurements have an important impact on the estimation of greenhouse gas (GHG) emissions from tree stems. Leaks are difficult to eradicate, even when using moldable cement, and the effects of concentration gradients are highly variable among gases. Based on this analysis, it is advisable to consider both phenomena. However, this implies determining  $k_L$  and  $C_b$ , which would require relatively long and tedious fieldwork. However, the determination of leaks and  $C_b$  as performed in the present study for explanatory purposes, is not mandatory. Data can be corrected using a simple analytical solution previously suggested (see below), although it is not commonly employed.

During the deployment of the chamber, it is observed that the concentration of  $C_C$  follows an asymptotic curve from an initial value  $C_{C,1}$  to a final value  $C_{C,2}$ , as the chamber reaches a steady-state condition and  $C_C$  becomes independent of time, as discussed by Welles et al. (2001) for soil CO<sub>2</sub> fluxes. This concentration trend can be described by a standard exponential equation that has been previously described by several authors (Hutchinson and Mosier, 1981; Pedersen et al., 2010; Pihlatie et al., 2013), similar to Eq. (6) (Eq. (14); specific details provided in Section S1.4).

$$C_{C,t} = C_{C,2} - (C_{C,2} - C_{C,1}) \cdot \exp(-k_t \cdot t) \quad (14)$$

where  $k_t$  is a kinetic constant which include all kinetic parameters of flux and leaks. Eq. (14) is numerically equivalent to Eq. (5), and can be used instead, without requiring the determination of  $k_L$  and  $C_b$ . Eq. (14) can be easily calibrated using any experimental  $C_C$  segment of the measurement series, using  $C_{C,2}$  and  $kt$  as adjustment parameters. After determining these parameters,  $F^*$  can be numerically determined, assuming  $C_C = C_A$ , as follows:

$$F^* = \left( \frac{dC_C}{dt} \right)_{t=0} \cdot \frac{V_T}{A_C} = k_t \cdot (C_{C,2} - C_A) \cdot \frac{V_T}{A_C} \quad (15)$$

Eq. (15), can be used in all cases, in substitution to Eq. (5), without requiring the identification of  $k_L$  and  $C_b$ .

Regarding the implications of our findings on the methods used for the determination of GHG emissions from tree stems, the dynamic closed chamber, i.e., continuous measurement of  $C_C$  in a closed chamber, is certainly the best method, as it would allow calibration of Eq. (14) with a large dataset. Discrete sampling in a closed chamber (static chamber method) could also be applied using this approach, but only if several samples are taken, to allow for Eq. (14) calibration.

In all cases, it is highly advisable to proceed to a good ventilation of the chamber prior to any measurement (Hutchinson and Livingston, 2001), in order to reach initial condition where  $C_C$  is as close as possible to  $C_A$ . The results obtained in the present work and the experimental strategy suggested indicate that the use of cement is not a requirement, although it would certainly provide better data quality in trees with severe roughness, i.e., reducing experimental noise. It is also advisable, in order to reduce noise at the onset of measurements, to use a pneumatic circuit as short as possible by placing the detector as close as possible to the chamber and to consider detectors with an internal pump flow rate as high as possible or even to consider an additional external recirculation pump. These recommendations are favorable to a short gas residence time within the pneumatic circuit, thus allowing a rapid equilibration between the chamber and the detector, producing better data quality at the onset of the measurement.

## 4. Conclusion

The effects of leaks and of concentration buildup severely impair flux measurements from tree stems and must be considered. After conducting a comprehensive literature review on CH<sub>4</sub> fluxes from tree stems over the past decade (69 items, Table S3), we observed that linear adjustment was consistently used in 75 % of the reports, while the choice of asymptotic-like data adjustment for flux calculation was only utilized in approximately 9 % of the cases. Consequently, the use of asymptotic fitting is not a common practice and remains limited, despite its ability to effectively handle data with or without significant curvature. Additionally, among the reports that employed linear adjustment, 42 % were based on discrete sampling methods, which complicates the accurate determination of potential data curvature over time caused by leaks and/or concentration buildup. By separately analyzing the impacts of leaks and concentration buildup, our study underscores the importance of both factors in accurately measuring CH<sub>4</sub> fluxes, and leaks having a dominant effect on CO<sub>2</sub> flux determination. Overall, this study serves as a reminder of the key findings highlighted in previous reports: the precise assessment of greenhouse gas emissions from tree stems necessitates the consideration of both leaks and concentration buildup, with the utilization of asymptotic data adjustment methods. These results, combined with the literature review, also suggest that our current estimates of GHG flux from tree stems are currently underestimated.

### CRedit authorship contribution statement

Julio A. Salas-Rabaza: Conceptualization, Funding and administrative management, Field/lab work México/Chile, Data processing, Writing - Original draft preparation; José Luis Andrade: Funding and administrative management, Revision - Original draft; Roberth Us-Santamaría: Fieldwork and logistic Mexico; Pablo Morales-Rico: Fieldwork Mexico, Gisela Mayora: Fieldwork Mexico; Francisco Javier Aguirre: Fieldwork and logistic Chile; Vicente Fecci-Machuca Field/lab work Chile; Eugenia M. Gade-Palma: Fieldwork Chile; Frederic Thalasso: Conceptualization, Funding and administrative management, Field/lab work México/Chile, Data processing, Writing - original draft preparation.

### Declaration of competing interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

### Data availability

Data will be made available on request.

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## Appendix A. Supplementary data

Supplementary data to this article can be found online at <https://doi.org/10.1016/j.scitotenv.2023.166358>.

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